

Physics Capabilities of a Large Liq-Ar Atmospheric detector

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(work with V. Barger, P. Ghoshal, S. Goswami, D. Marfatia, S.
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arXiv 1203.6012



PXPS Meeting, Fermilab, June 18, 2012

Atmospheric detectors... Some strengths and limitations

Atmospheric detectors collect neutrino fluxes which are uncertain by 10 - 20% in absolute value (Flavour ratios less uncertain)

Fluxes are sharply dropping functions of energy,

$$\frac{d\phi_\nu}{dE} \sim E^{-\gamma} \text{ and } \gamma \simeq 3 \text{ for } \nu_\mu \text{ and } \simeq 3.5 \text{ for } \nu_e$$

At energies where flux is significant (below $\sim 2 \text{ GeV}$), neutrino-nucleon cross-section has significant uncertainties

Also, at energies where flux is significant (below $\sim 2 \text{ GeV}$), produced charged lepton may have significantly different direction from incoming neutrino

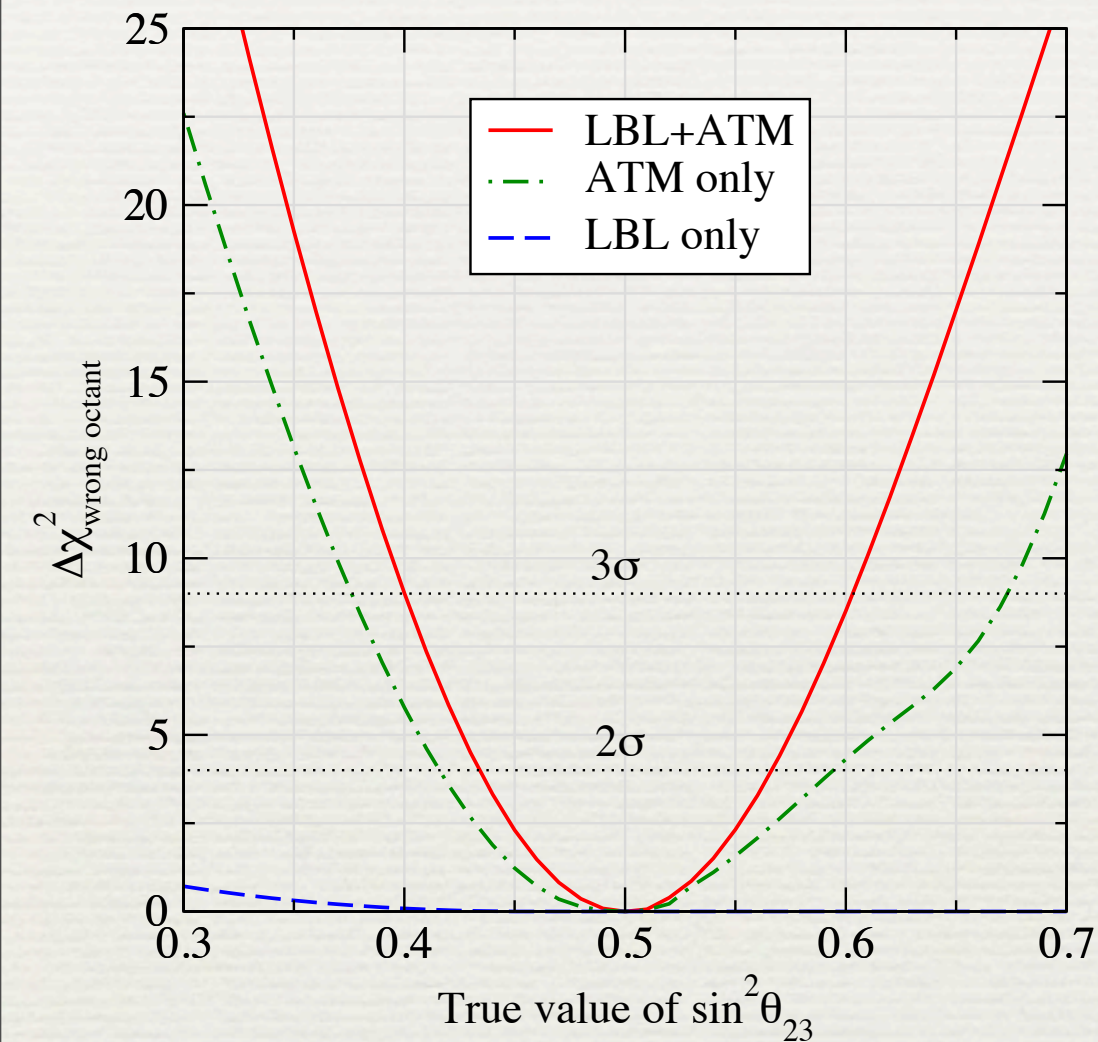
Atmospheric detectors... Some strengths and limitations

Thus there is an inherent lack of precision in L and in E on an event-by-event basis in atmospheric detectors, which translates into uncertainties in measurements of Δm^2 and $\sin^2 2\theta_{23}$

Unlike LBL experiments, atmospheric detectors, however, tap into a very broad band in L (20 km to 12500 km) and E (100 MeV to 10 TeV).

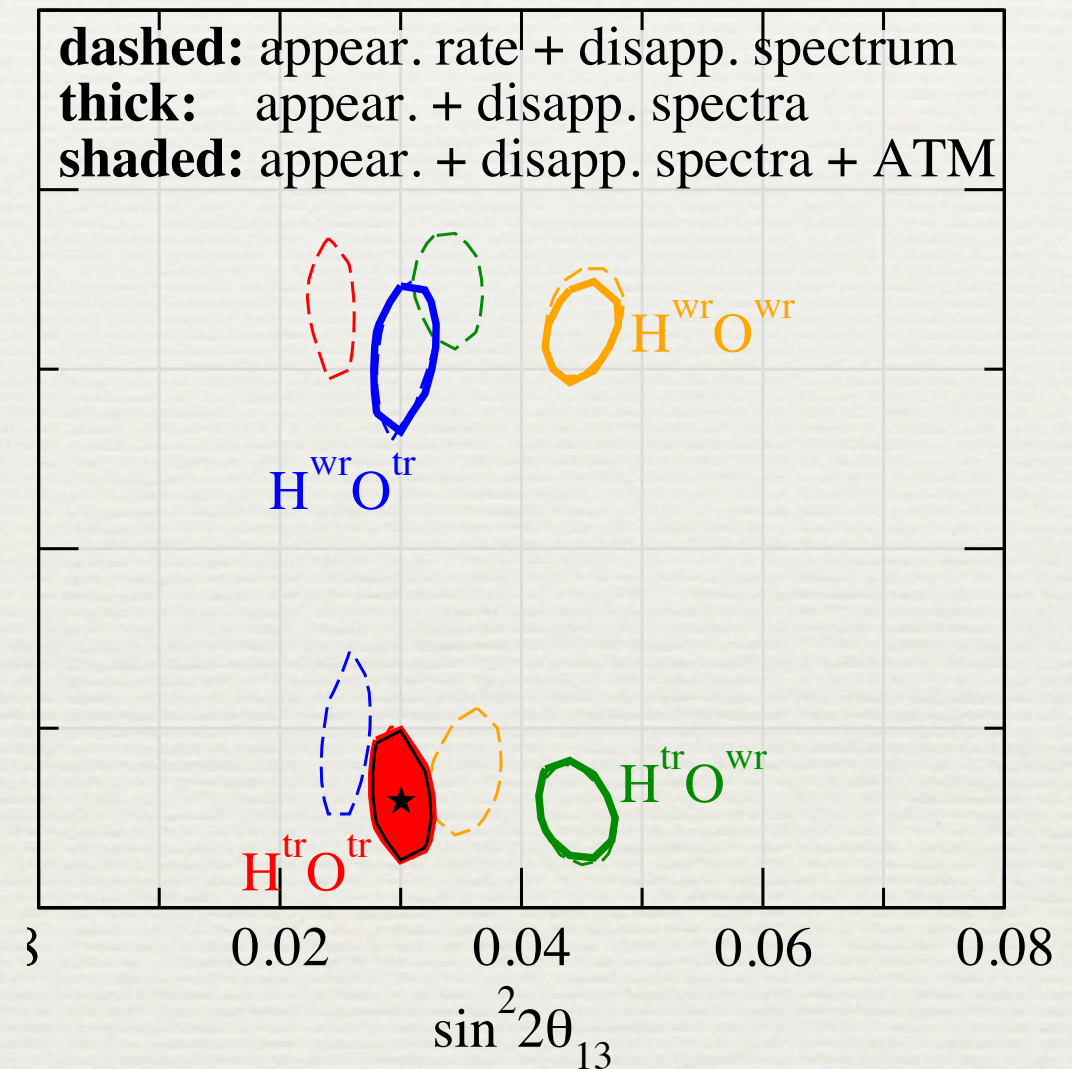
Allows resolution of parameter degeneracies which are a feature of LBL set-ups.

Earlier studies have shown the power of combining LBL + ATM data to maximize the physics output.....



Huber, Maltoni and Schwetz
hep-ph 0501037

T2K-II + HK, 295 km

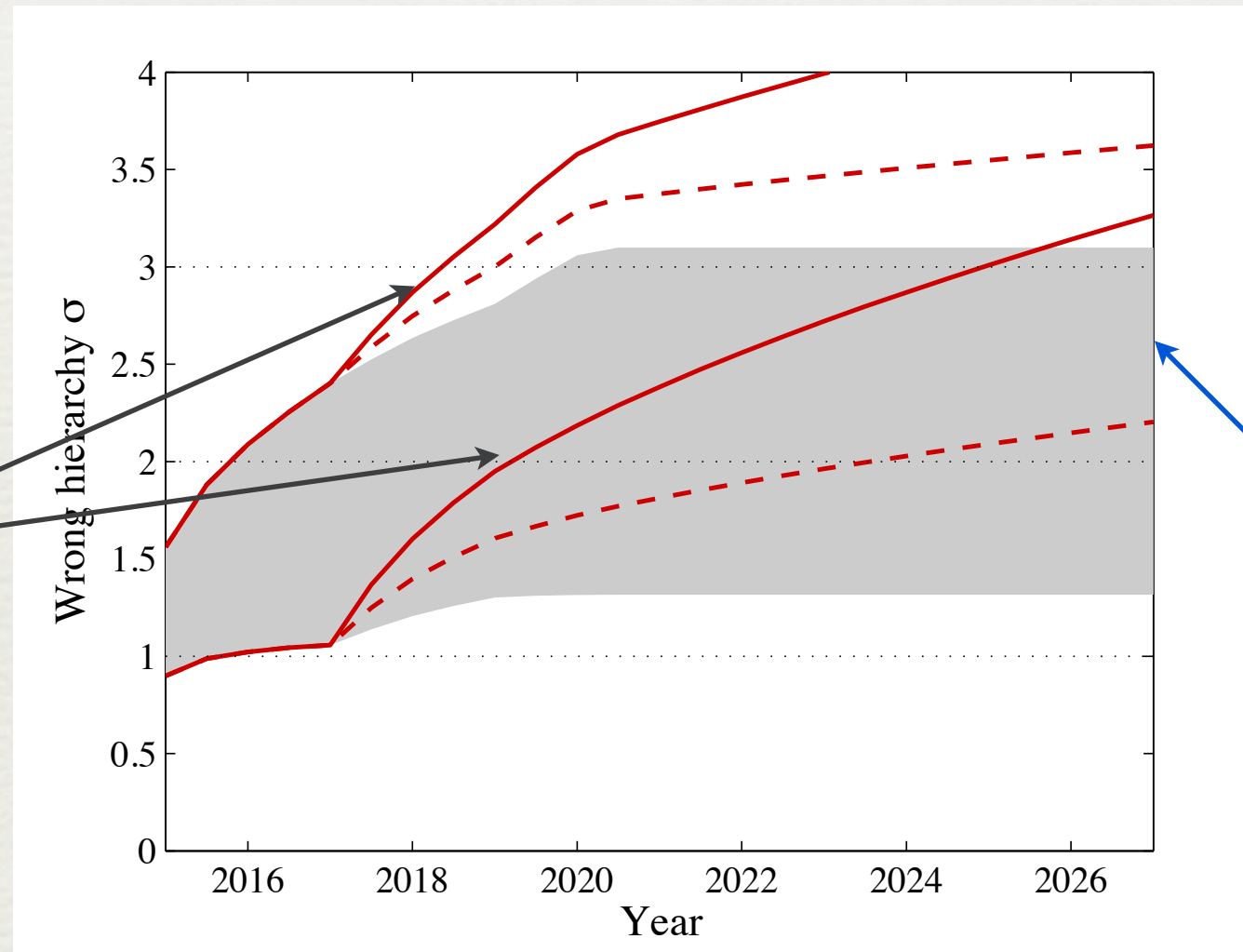


Campagne, Maltoni, Mezzetto and Schwetz
hep-ph 0603172

SPS + MEMPHYS, 130 km

Earlier studies have shown the power of combining LBL + ATM data to maximize the physics output.....

INO + T2K + NO ν A

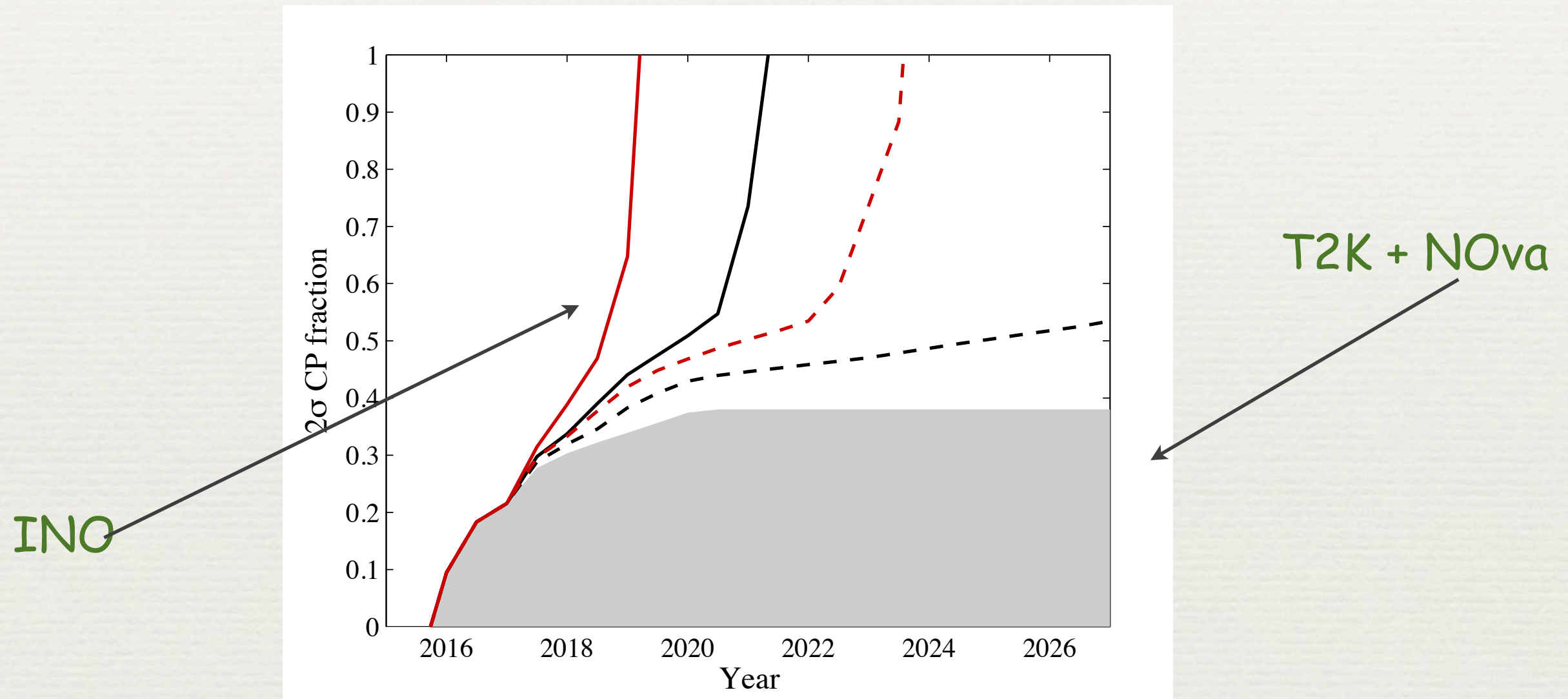


T2K + NO ν A

Blennow and Schwetz
arXiv 1203.3388

INO + T2K + NO ν A, hierarchy sensitivity,

Earlier studies have shown the power of combining LBL + ATM data to maximize the physics output.....



Large θ_{13}Some consequences

Recent result from Daya Bay $\sin^2 2\theta_{13} = 0.092 \pm 0.016(stat) \pm 0.005(syst)$

Recent result from RENO $\sin^2 2\theta_{13} = 0.113 \pm 0.013(stat) \pm 0.019(syst)$

Measurements significantly impact the planning of future neutrino facilities.

Large θ_{13}Consequences for Atmospheric Detectors

Importance of future large mass atmospheric detectors has increased, e.g INO and the case for a Liq Ar detector, because large θ_{13} allows them to tap into matter effects better, and provides a shot at determining the mass hierarchy and octant of θ_{23}

The case for combining super-beam facilities with large mass atmospheric detectors analytically and both analytically and physically (LBNE with large Liq Ar) has become stronger.

Large θ_{13}Some consequences

$$P_{\alpha\beta} = P_{\alpha\beta}^{(0)} + \alpha P_{\alpha\beta}^{(1)} + \mathcal{O}(\alpha^2)$$

$$P_{e\mu}^{(0)} = s_{23}^2 \frac{\sin^2 2\theta_{13}}{C_{13}^2} \sin^2 C_{13} \Delta,$$

$$C_{13} \equiv \sqrt{\sin^2 2\theta_{13} + (A - \cos 2\theta_{13})^2}.$$

$$P_{e\mu}^{(1)} = -2s_{12}^2 s_{23}^2 \frac{\sin^2 2\theta_{13}}{C_{13}^2} \sin C_{13} \Delta$$

$$\Delta \equiv \frac{\Delta m_{31}^2 L}{4E},$$

$$\times \left[\Delta \frac{\cos C_{13} \Delta}{C_{13}} (1 - A \cos 2\theta_{13}) - A \frac{\sin C_{13} \Delta}{C_{13}} \frac{\cos 2\theta_{13} - A}{C_{13}} \right]$$

$$A \equiv \frac{2EV}{\Delta m_{31}^2} = \frac{VL}{2\Delta}.$$

Akhmedov et al, hep-ph 0402.175

$$+ s_{13} \sin 2\theta_{12} \sin 2\theta_{23} \frac{\sin C_{13} \Delta}{A C_{13}^2} \left\{ \sin \delta_{\text{CP}} [\cos C_{13} \Delta - \cos(1 + A)\Delta] C_{13} \right.$$

$$\left. + \cos \delta_{\text{CP}} [C_{13} \sin(1 + A)\Delta - (1 - A \cos 2\theta_{13}) \sin C_{13} \Delta] \right\},$$

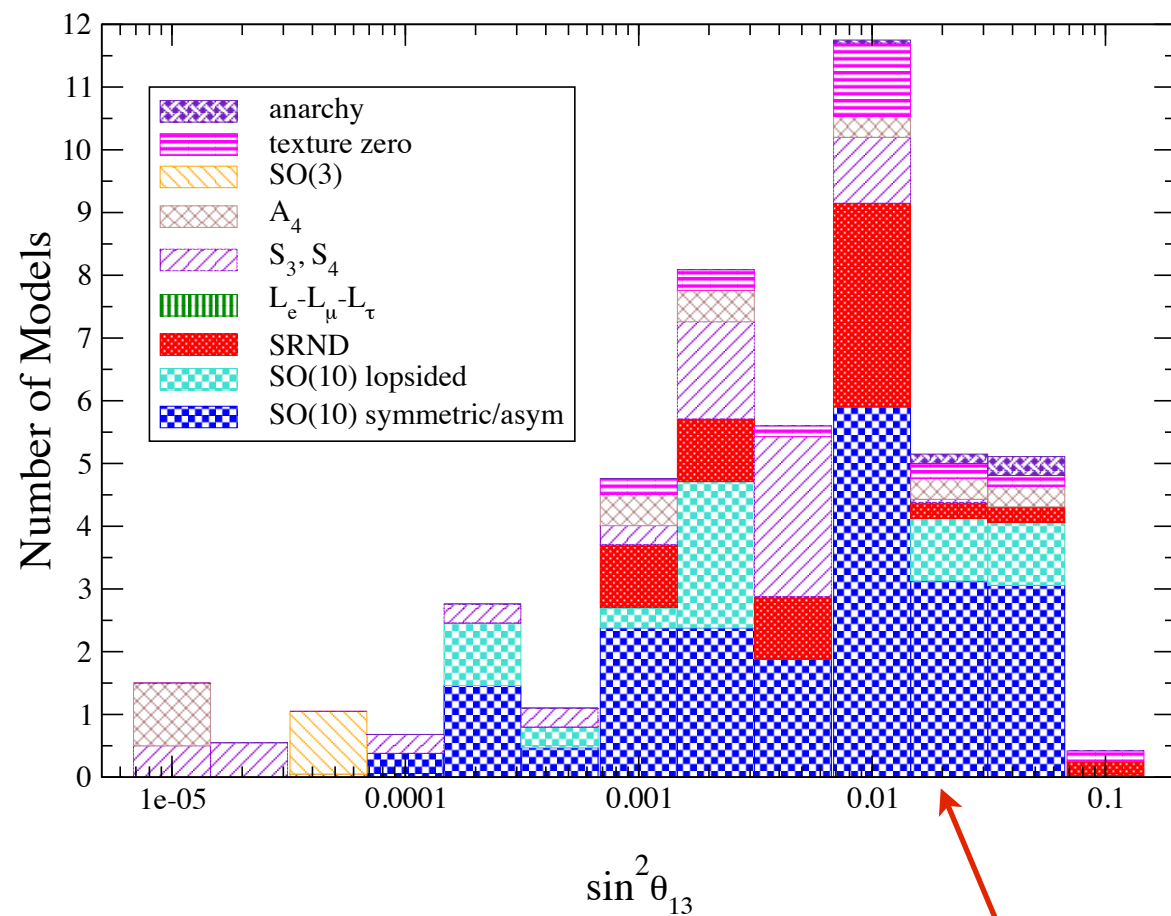
$$A_{CP} = \frac{P_{\nu_e \rightarrow \nu_\mu} - P_{\bar{\nu}_e \rightarrow \bar{\nu}_\mu}}{P_{\nu_e \rightarrow \nu_\mu} + P_{\bar{\nu}_e \rightarrow \bar{\nu}_\mu}} \sim \frac{1}{\sin \theta_{13}}$$

Large θ_{13} does not automatically mean enhanced CP measurement.

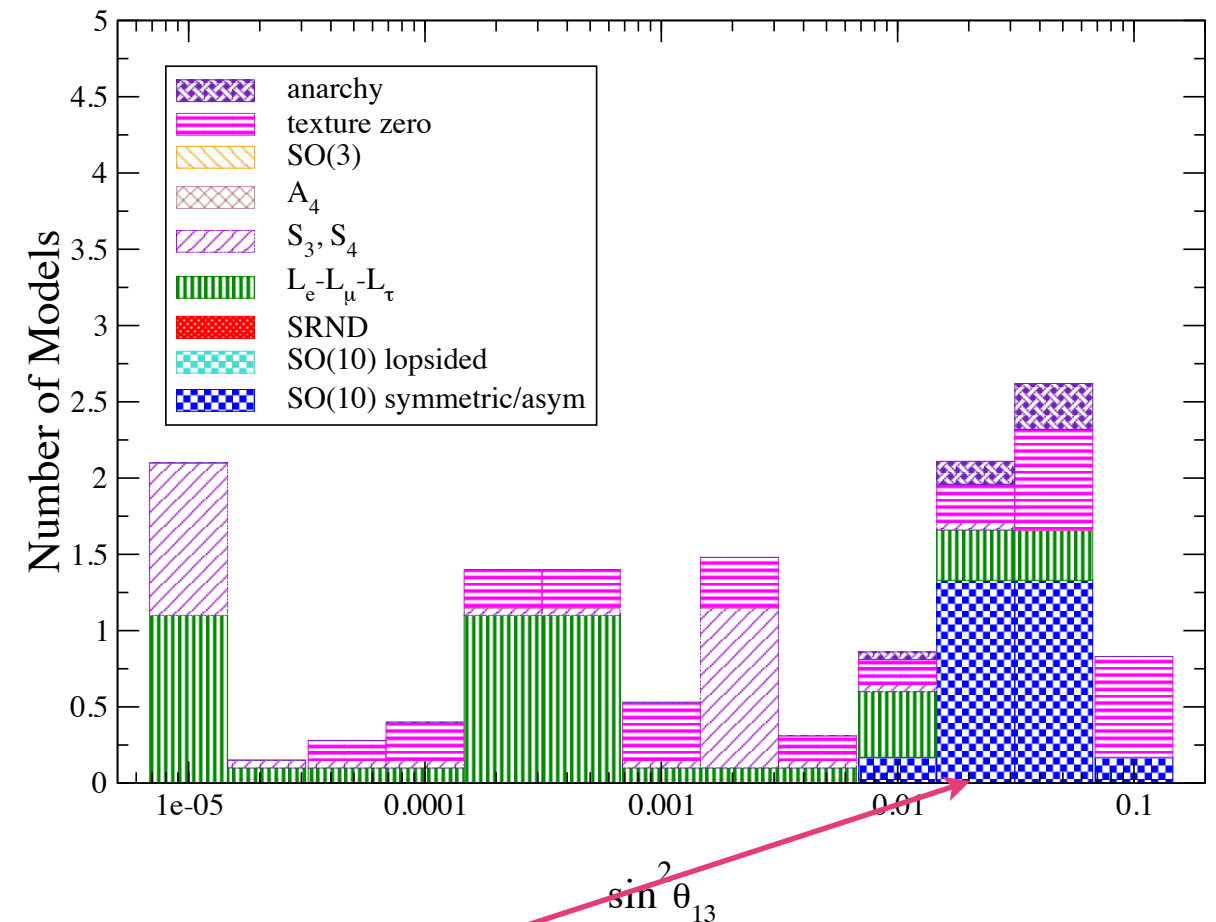
Overall CP sensitivity depends on baseline and type of expt

Importance of the mass hierarchy.....

Models with Normal Hierarchy



Models with Inverted Hierarchy

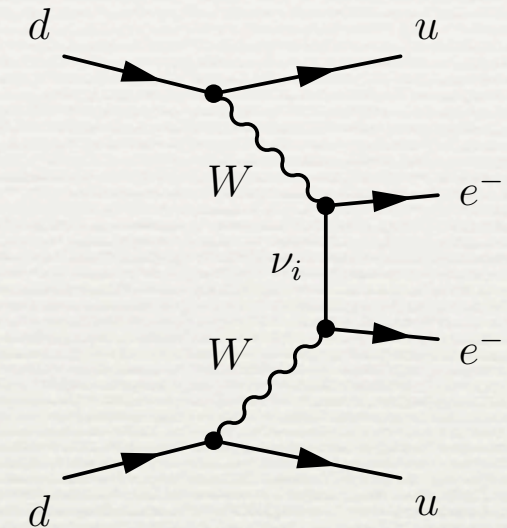
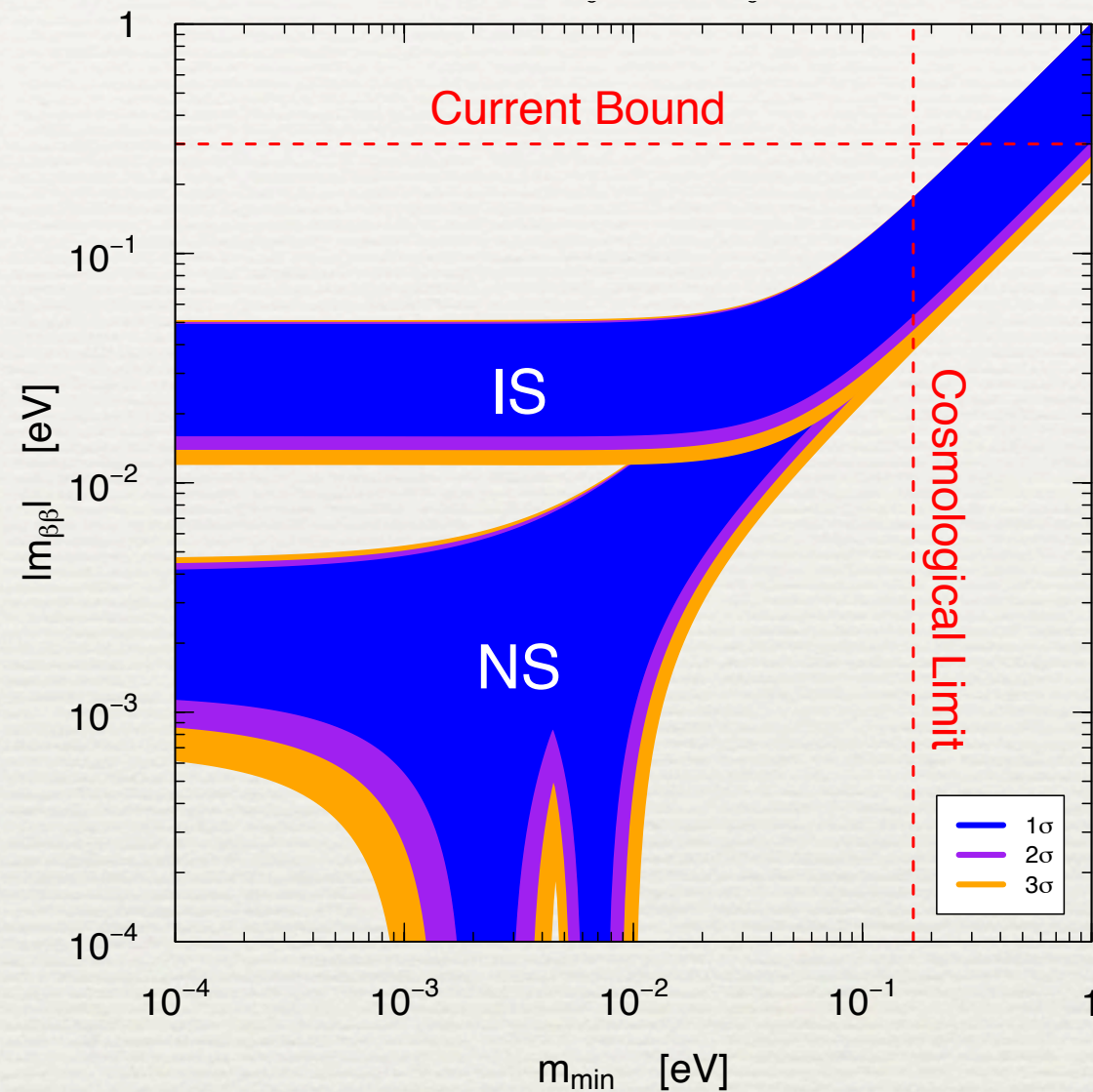


$\sin^2 2\theta_{13} = 0.1$

Albright and Chen, hep-ph 0608137

A very useful discriminator of BSM models.....

Importance of the mass hierarchy.....

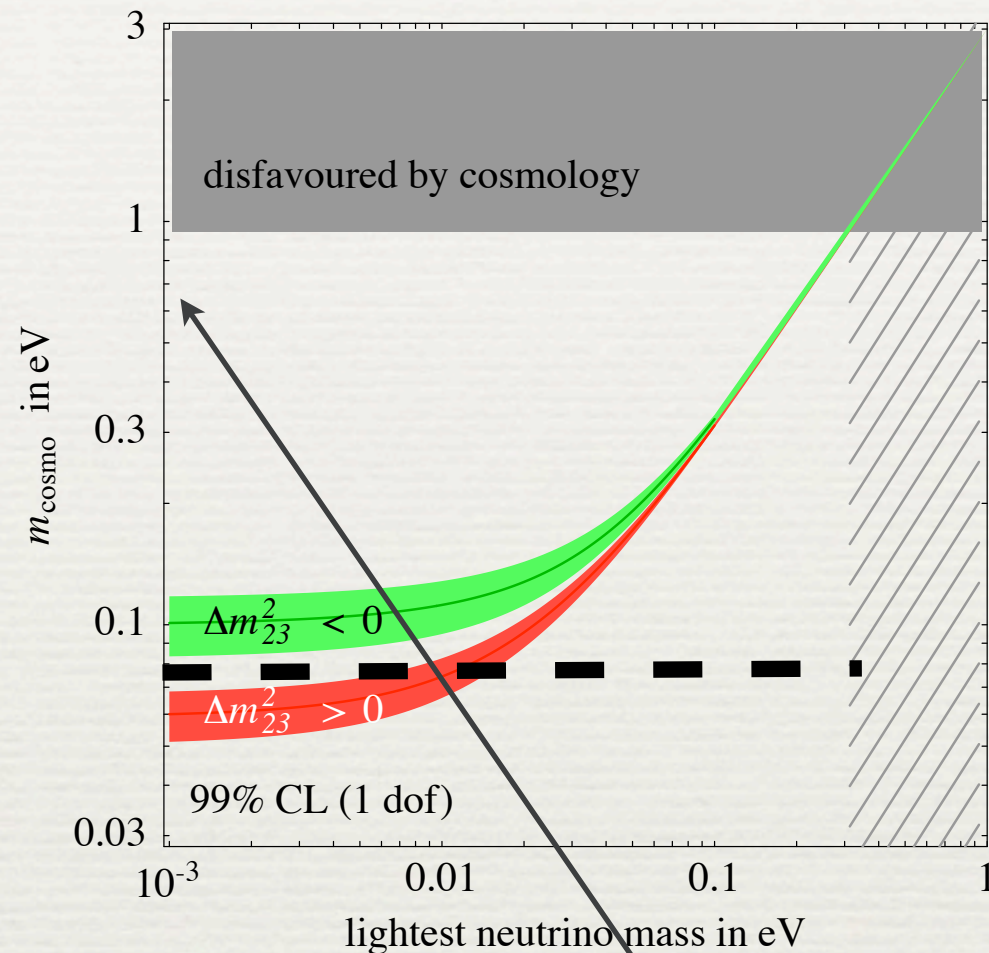


Bilenky and Giunti, 1203.5250

$$m_{\beta\beta} = \sum_i U_{ei}^2 m_i.$$

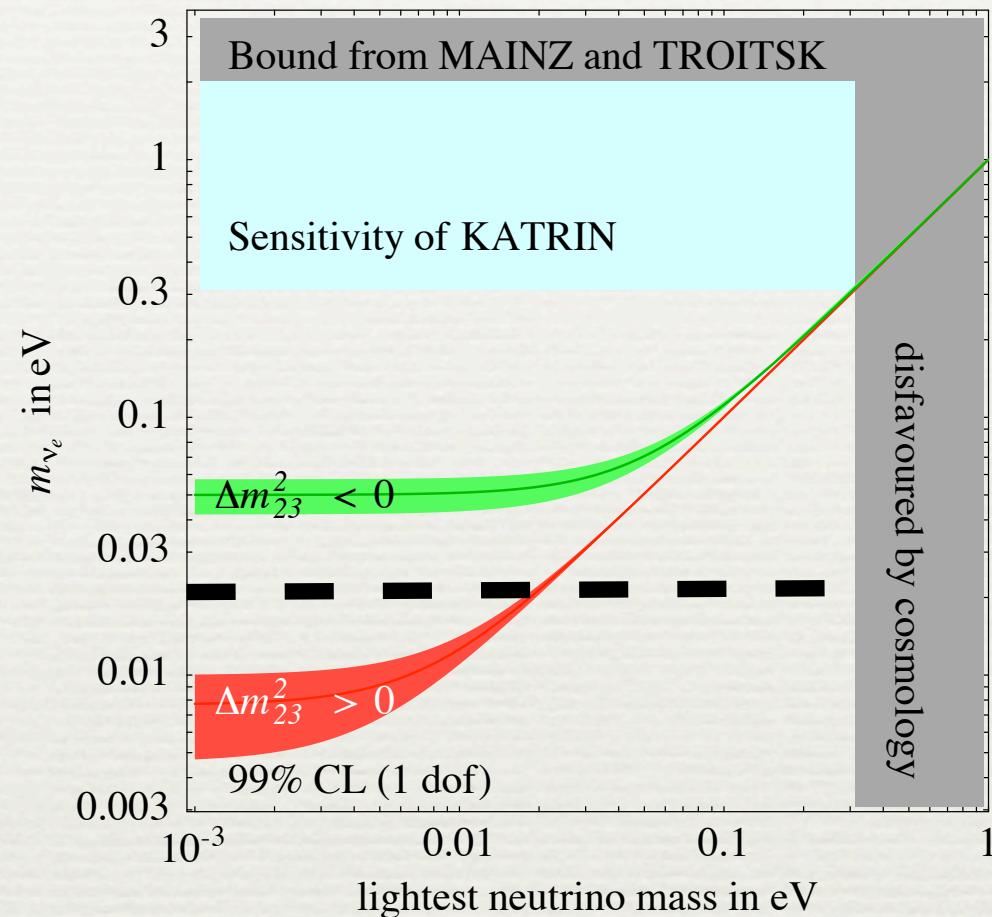
Width of $0\nu\beta\beta$ -decay depends on $|m_{\beta\beta}|^2$, the effective majorana mass squared.

Importance of the mass hierarchy.....



WMAP/SDSS/Hubble current bound
(conservative)

R. De Putter et al, 1201.1909, M Moresco et al,
1201.6658

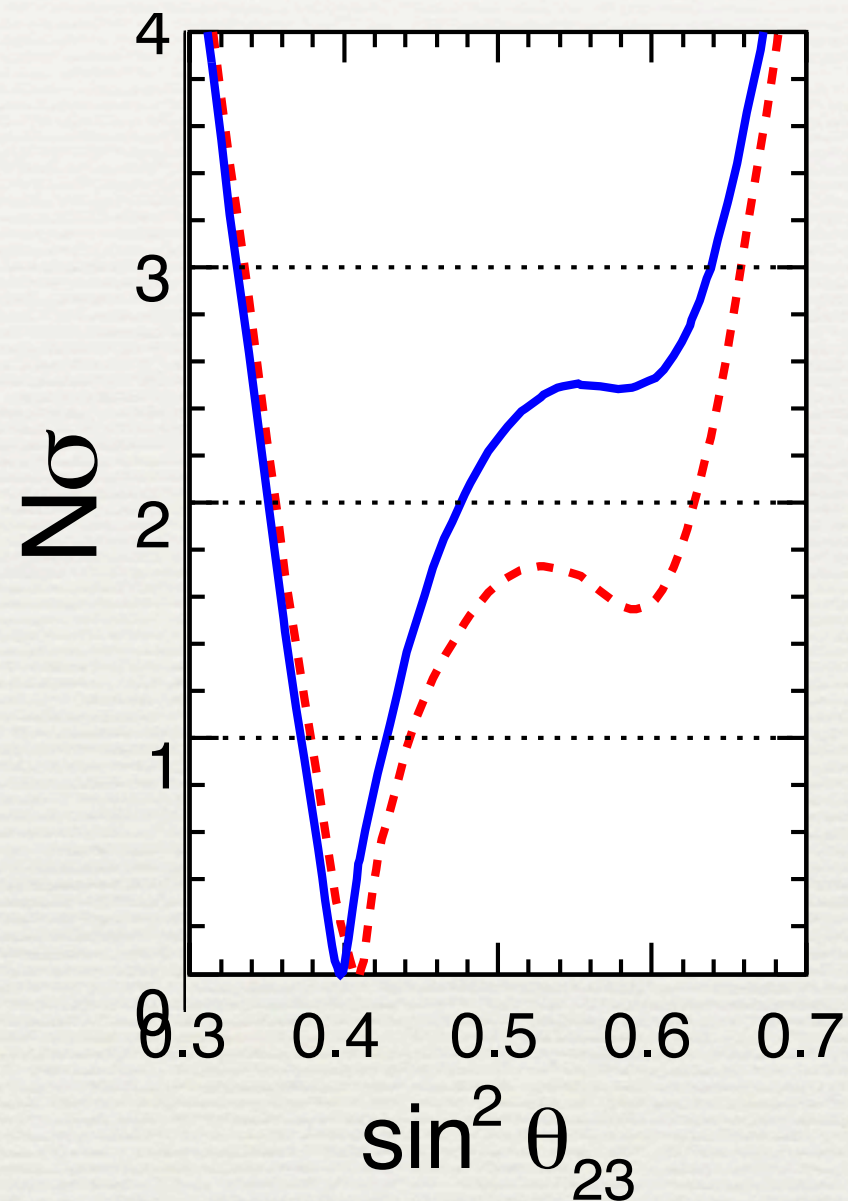


Strumia and Vissani, hep-ph 0503246
Parke, FNAL CONF 06-248-T

$$m_{\beta} = \sqrt{\sum_i |U_{ei}|^2 m_i^2}.$$

The link between hierarchy and sensitivity requirements of
experiments.....

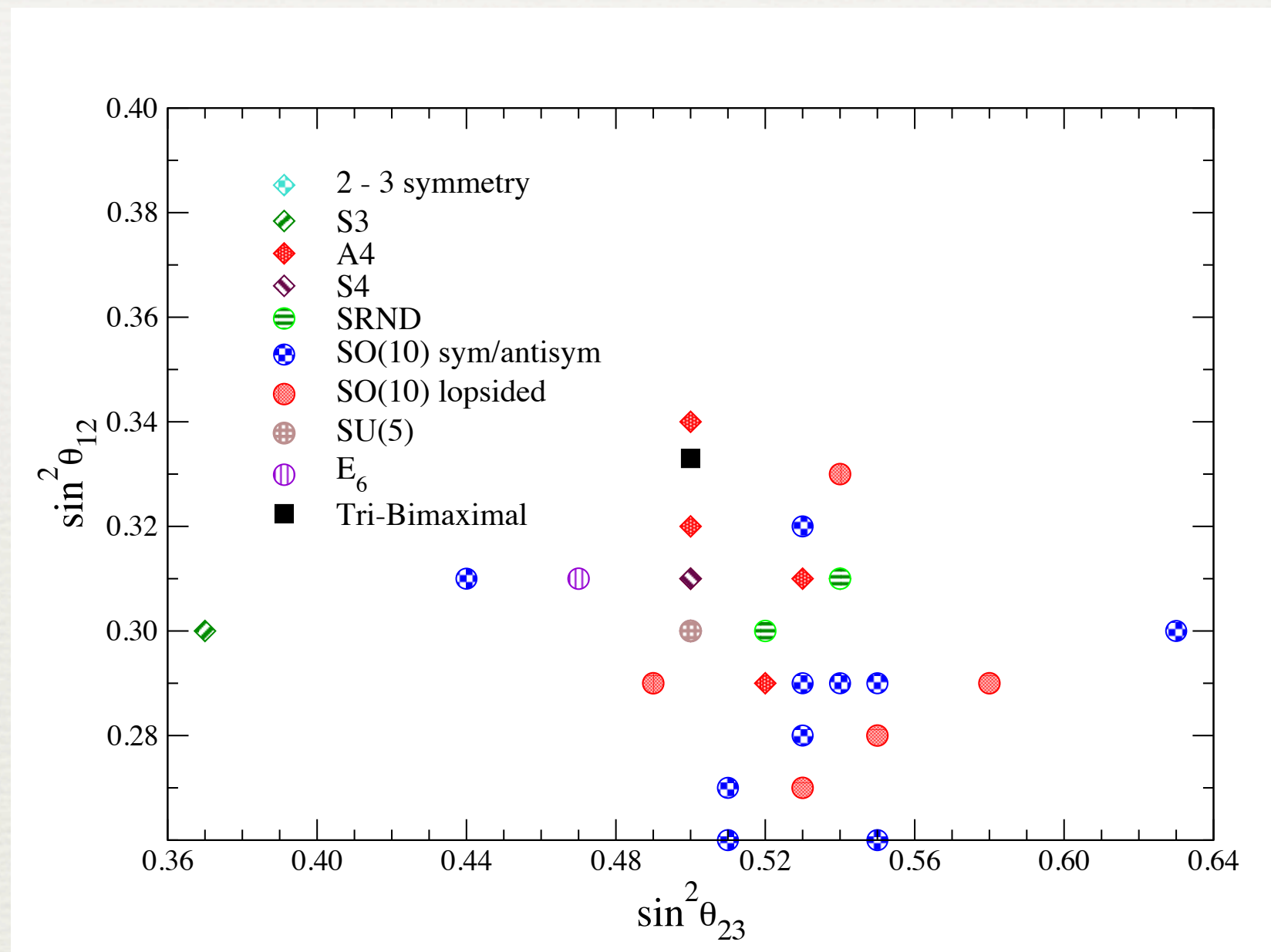
Octant of θ_{23}



Fogli et al, arXiv 1205.5254

Present Global data exhibits ~ 2 sigma preference for first octant, indicating a significant deviation from maximal mixing.

Octant of θ_{23}



Albright, arXiv 0905.0146

Discrimination among models using the value of θ_{23}

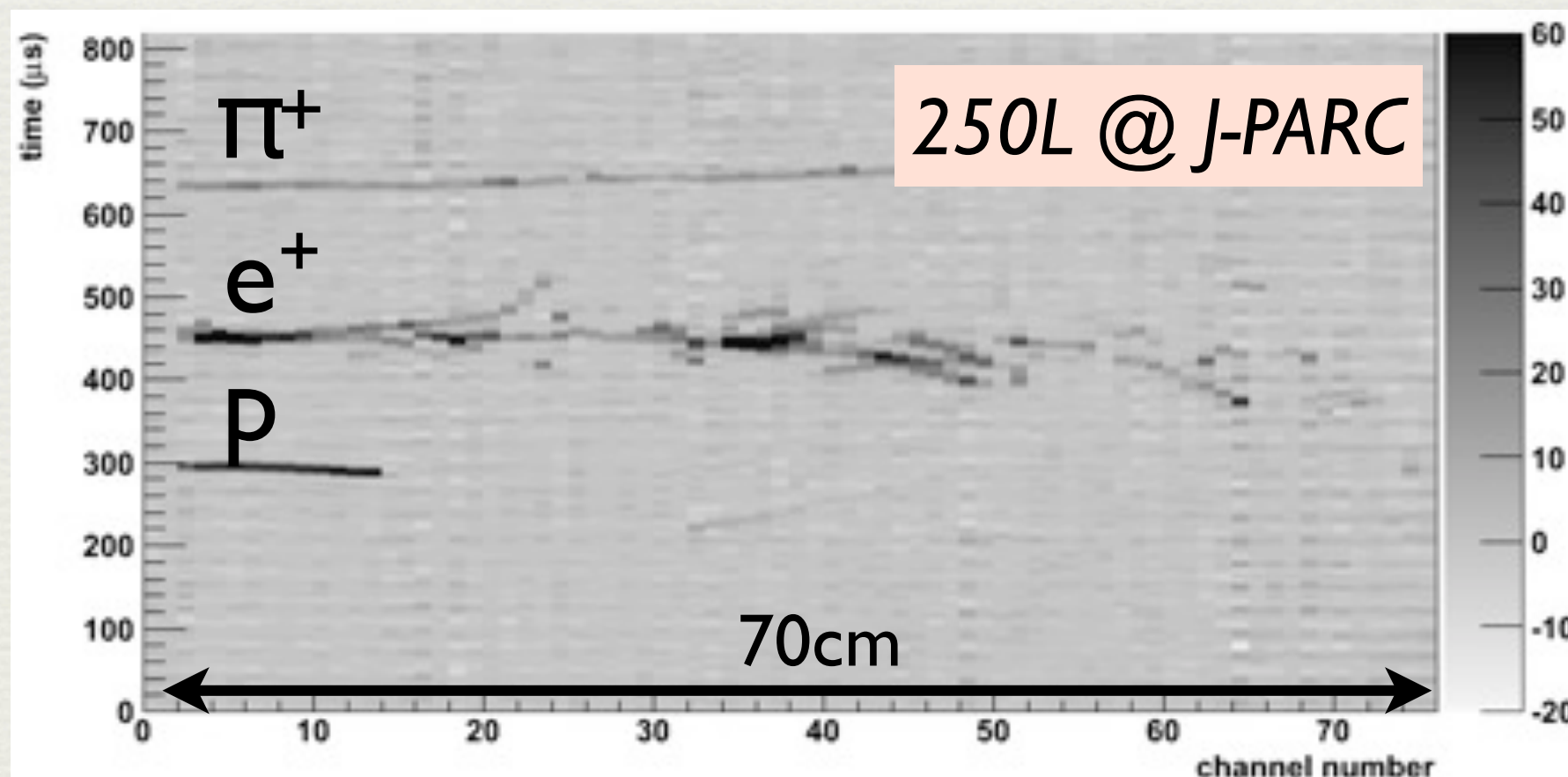
Liquid Argon Detectors.....

A Rubbia

Particle ionization tracks can be drifted over several meters leading to homogeneous and high accuracy imaging over large volumes.

Bubble chamber-like 3D imaging and mm scale spatial resolution.....

Particle id via dE/dx and/or decay product id



Liquid Argon Detectors.....magnetization

Studies done for Liq-AR TPC's for neutrino factories have examined the feasibility and advantages of magnetization of such detectors

For long tracks 4 m (e.g muons) a field of 0.1 T can provide >3 sigma charge id in this medium. Thus low threshold possible (800 MeV for muons)

Early showering of electrons makes their charge id difficult, but with 1 T field can achieve id for energies 1-5 GeV with 20% efficiency.

A Rubbia arXiv 0908.1286, A Rubbia et al hep-ph/0106088, A Bueno et al hep-ph/0005007, A Bueno et al hep-ph/0010308.....

LBNE Reconfiguration options.....Phase I

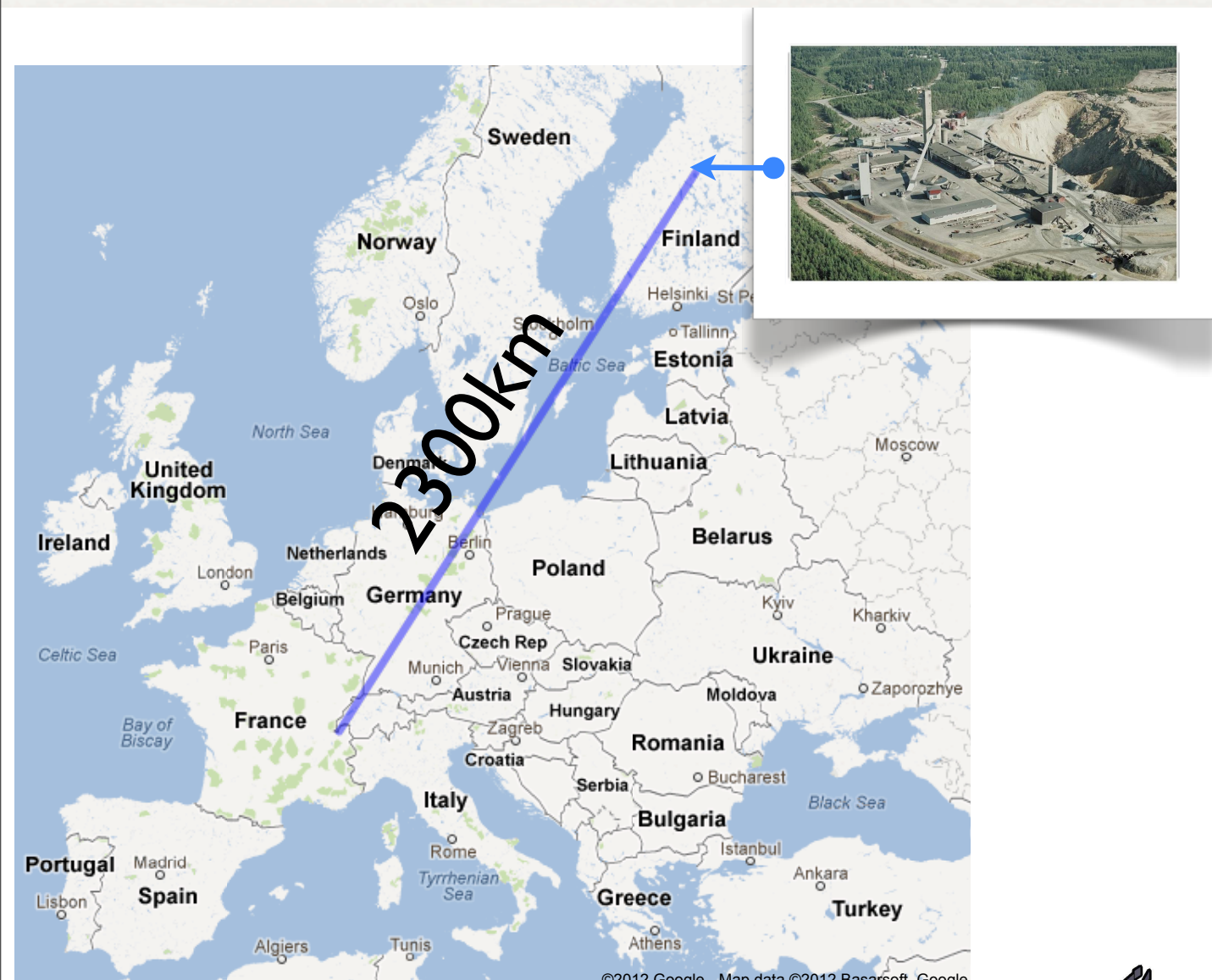
New 700 kW beam line to Homestake, no Near Detector, 10 kton Far Detector (would be on surface unless additional funding obtained)

Use existing NuMI beam and build a 15 kton detector underground at the MINOS site in the Soudan Mine.

Use existing NuMI beam and build a 30 kton detector on the surface at the NOVA site at Ash River.

Decision soon!

Liquid Argon Detectors..... Future



GLACIER

GLACIER/LAGUNA....Pyhasalami, Finland site under serious consideration, 20 kT (initial, w/subsequent increments) underground at 4000 m.w.e, wideband SPSbeam, 2300 km + Atmospheric neutrinos

Liquid Argon Detectors.....

A Bueno et al, hep-0701101
A Rubbia

Good calorimetric and angular resolution

$$\begin{aligned}\sigma_{E_e} &= \sigma_{E_\mu} = 0.01, \\ \sigma_{E_{had}} &= \sqrt{(0.15)^2/E_{had} + (0.03)^2}, \\ \sigma_{\theta_e} &= 0.03 \text{ radians} = 1.72^\circ, \\ \sigma_{\theta_\mu} &= \sigma_{\theta_{had}} = 0.04 \text{ radians} = 2.29^\circ,\end{aligned}$$

Leads to

$$\begin{aligned}\sigma_{E_\nu} &= \sqrt{(0.01)^2 + (0.15)^2/(yE_\nu) + (0.03)^2}, \\ \sigma_{\theta_{\nu e}} &= 2.8^\circ, \quad \sigma_{\theta_{\nu\mu}} = 3.2^\circ\end{aligned}$$

with

$$E_{had} = yE_\nu$$

Inputs into calculation.....

Parameter ranges for marginalization.....

- θ_{23} from $38^\circ - 52^\circ$ ($\sin^2 \theta_{23} = 0.38 - 0.62$),
- $|\Delta m_{31}^2|$ from $(2.05 - 2.75) \times 10^{-3} \text{ eV}^2$,
- θ_{13} from $5.5^\circ - 11.0^\circ$ ($\sin^2 2\theta_{13} = 0.04 - 0.14$),
- δ_{CP} from $0 - 2\pi$.

True values.....

$$(|\Delta m_{31}^2|)_{tr} = 2.4 \times 10^{-3}$$

$$(\delta_{CP})_{tr} = 0, (\sin^2 \theta_{23})_{tr} = 0.4, 0.5 \text{ and } 0.6$$

Fixed values.....

Δm_{21}^2 and θ_{12} are fixed at their best-fit

Errors, Uncertainties and priors.....

Flux normalization error 20%

Flux tilt factor, 5%

Overall cross-section uncertainty 10%

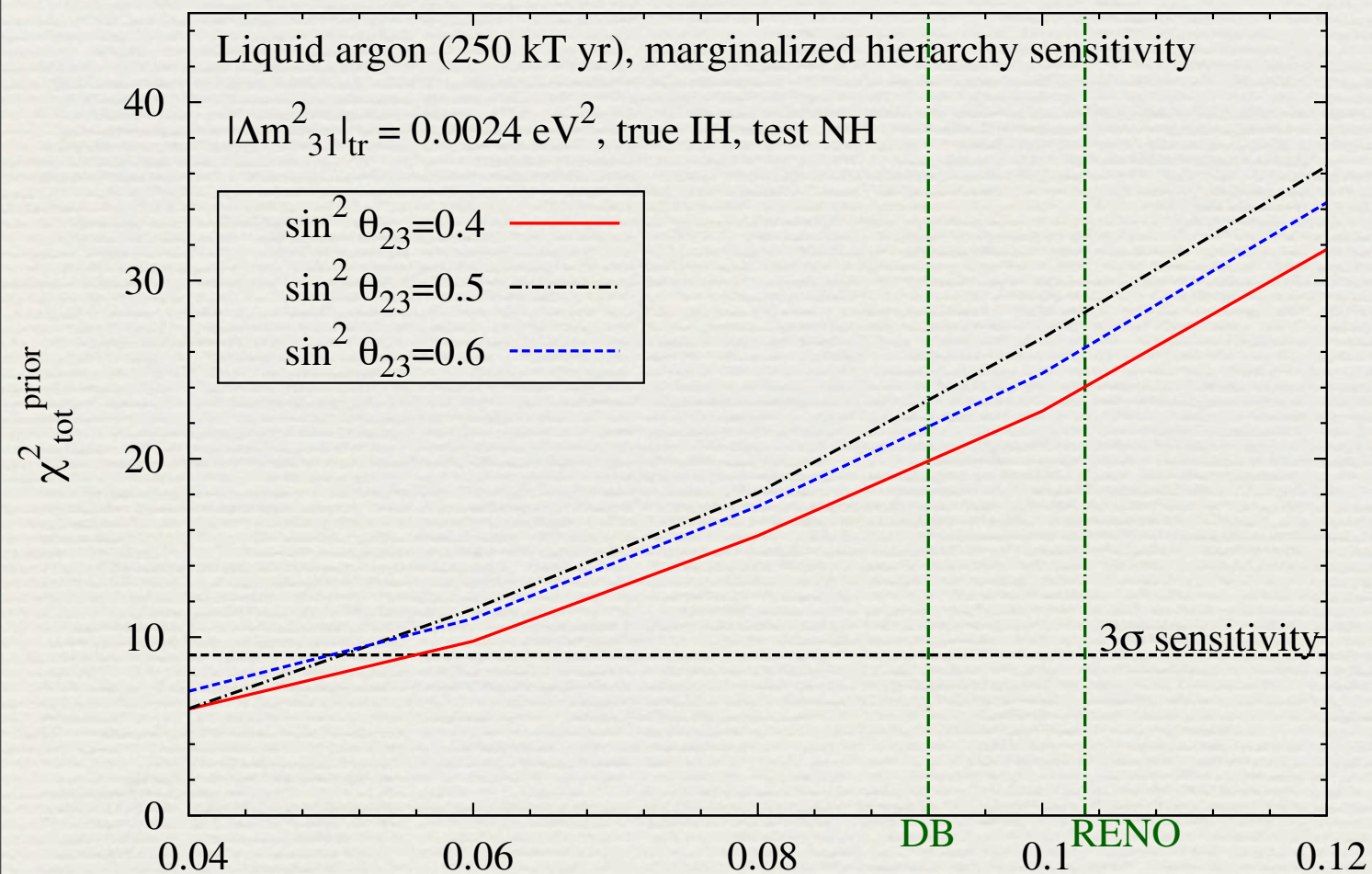
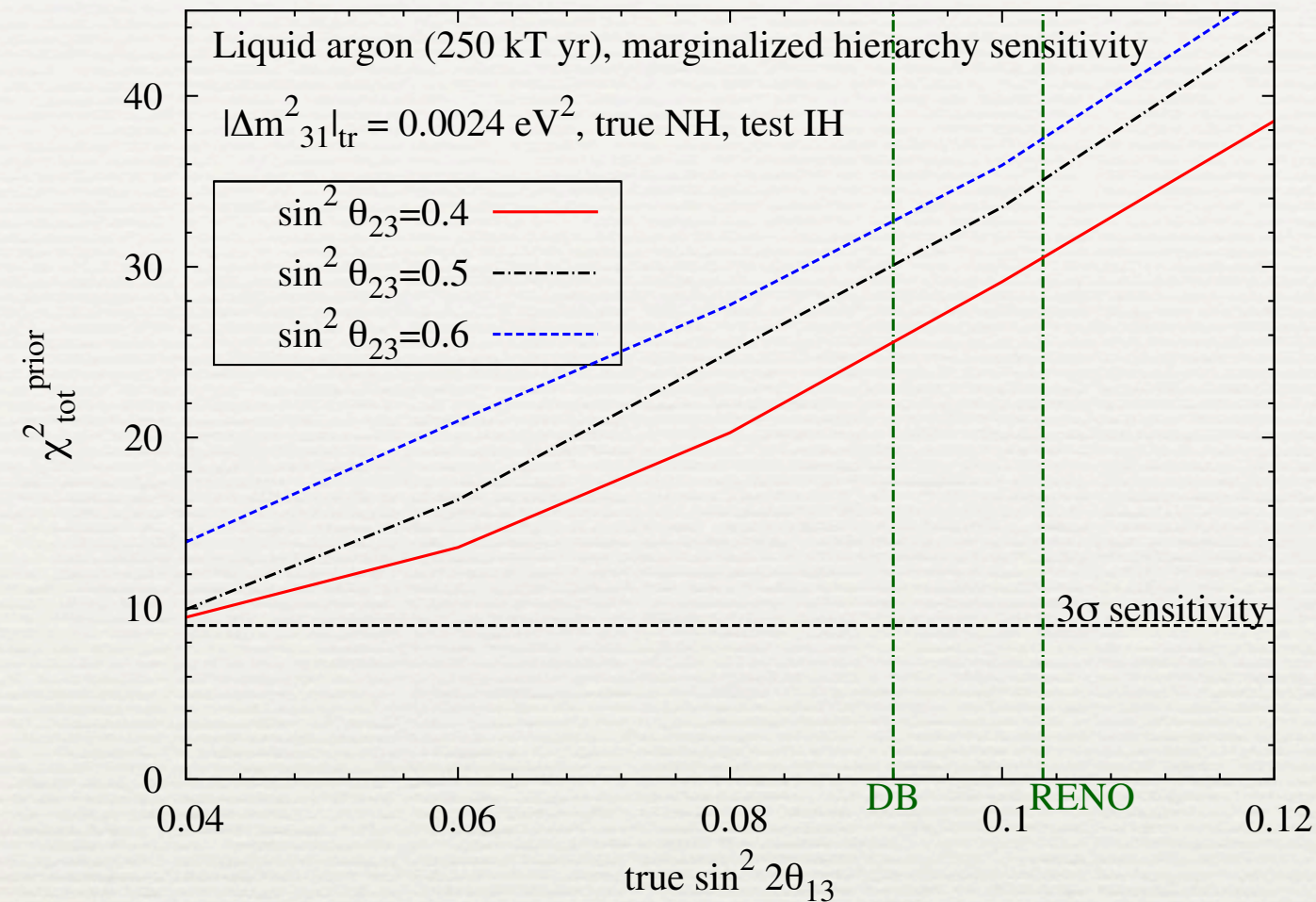
Overall systematic uncertainty 5%

Priors with 1σ ranges: $\sigma(|\Delta m_{31}^2|) = 0.05|\Delta m_{31}^2|$,
 $\sigma(\sin^2 2\theta_{13}) = 0.02$, and $\sigma(\sin^2 2\theta_{23}) = 0.02 \sin^2 2\theta_{23}$.

Minimize

$$\chi_\mu^2 + \chi_e^2 + \chi_{prior}^2.$$

Results.....Hierarchy



Sensitivity strongly depends on matter effects, the detection of which relies on charge id

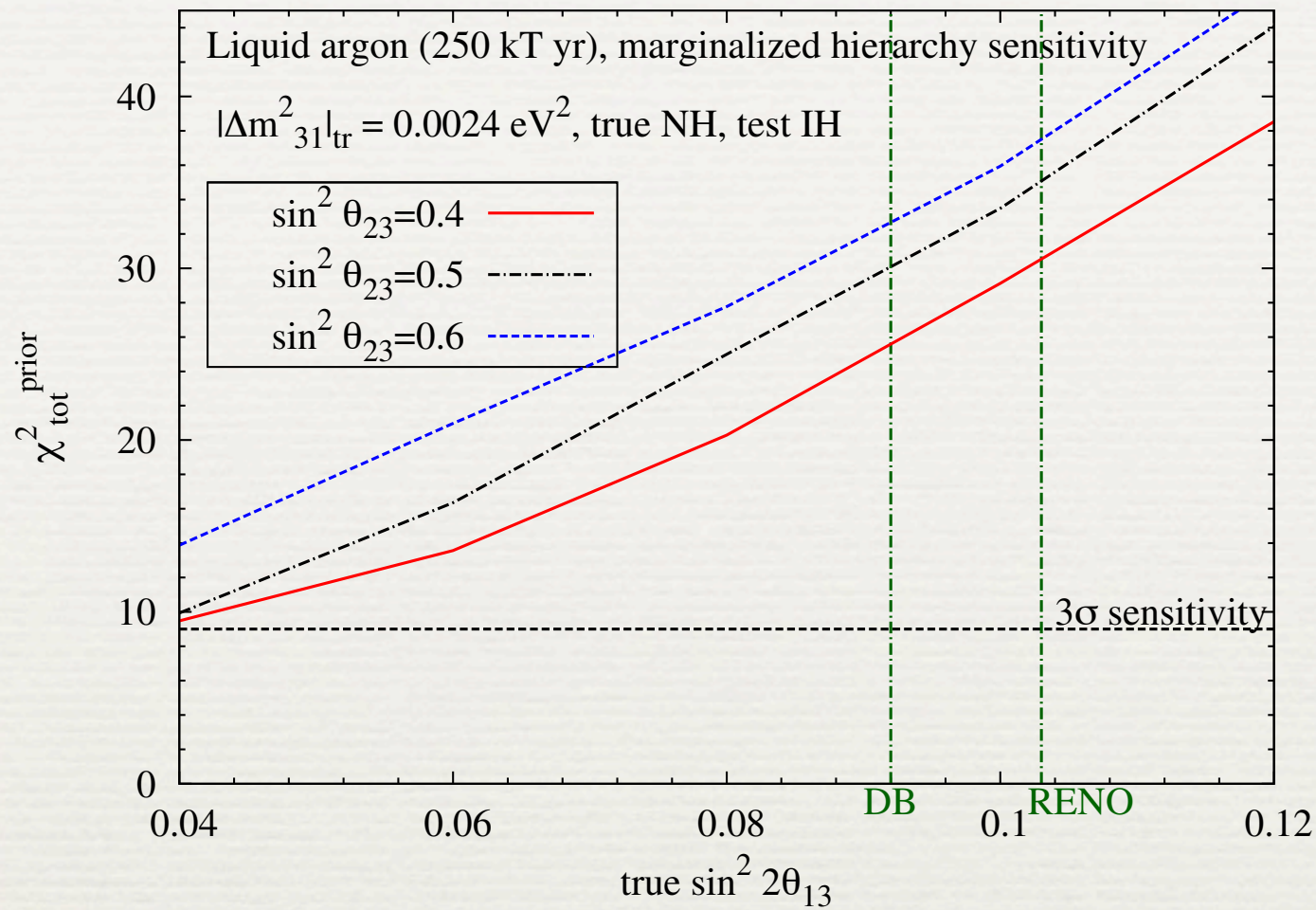
Muons make a larger contribution to sensitivity compared to electrons.

Superior angular resolution of the detector plays an important role

Priors and precision in θ_{13} important.

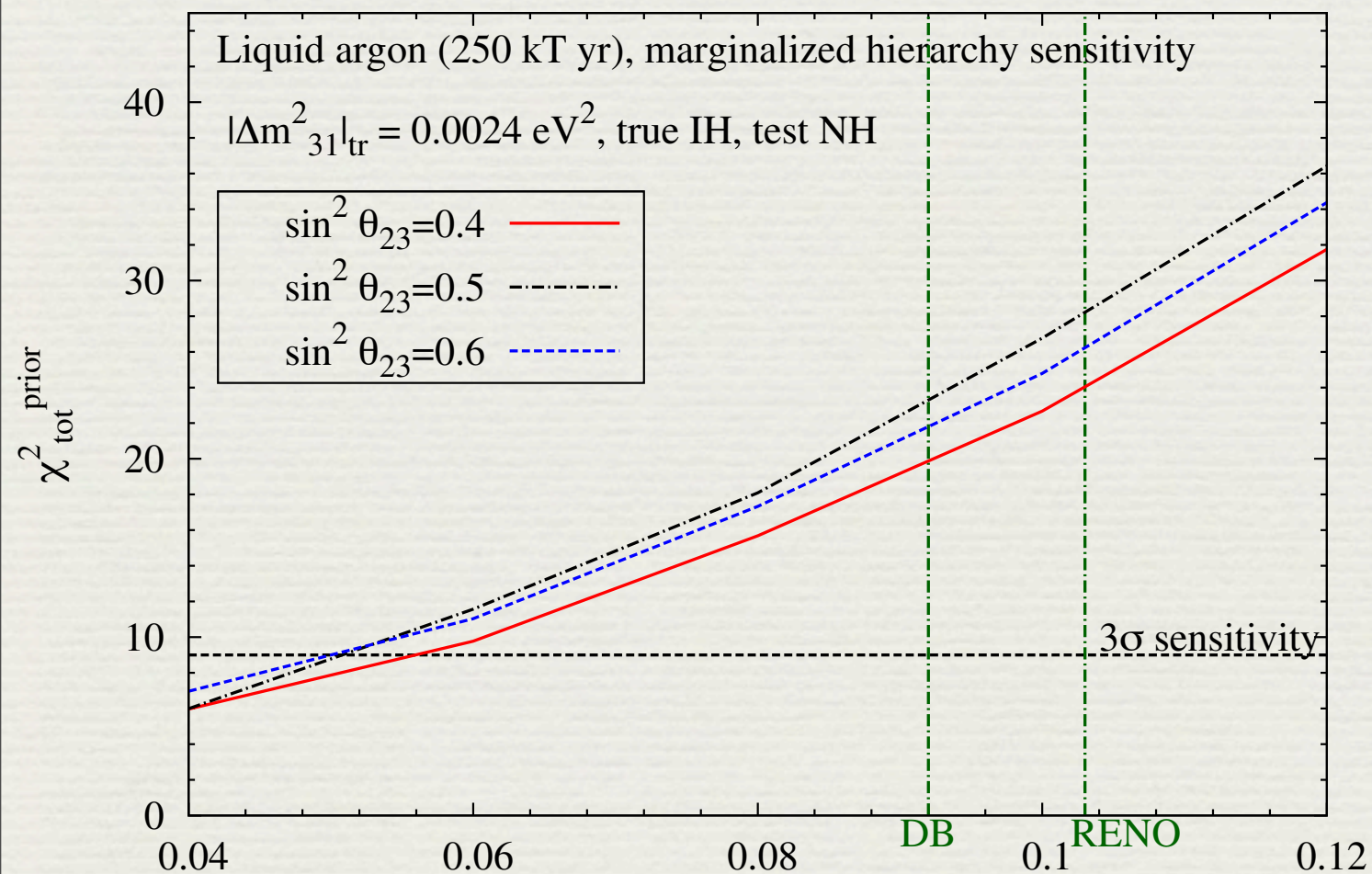
$\sigma(\sin^2 2\theta_{13})$	χ^2_{NH}	χ^2_{IH}
No Prior	25.6	11.7
0.02	33.5	26.8
0.01	35.8	36.8

Results.....Hierarchy

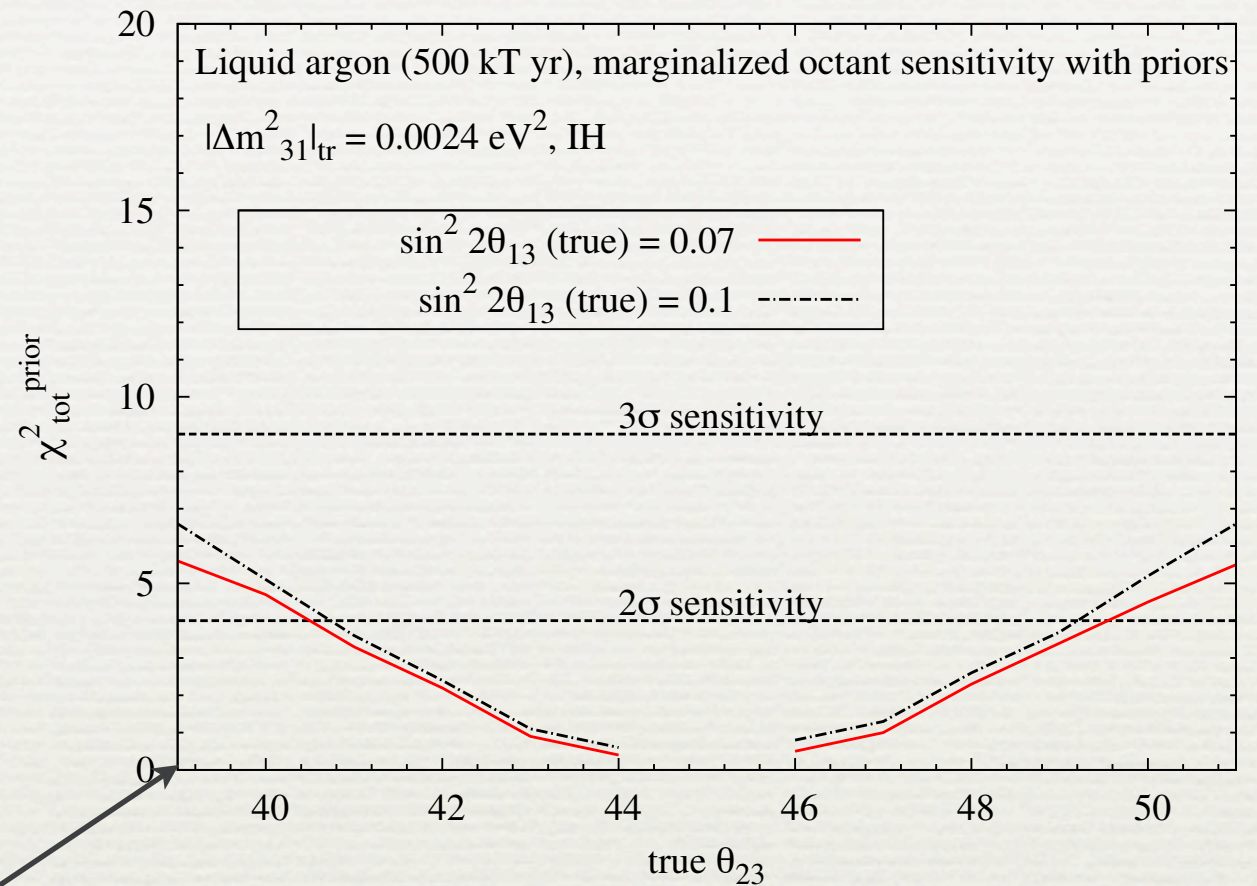
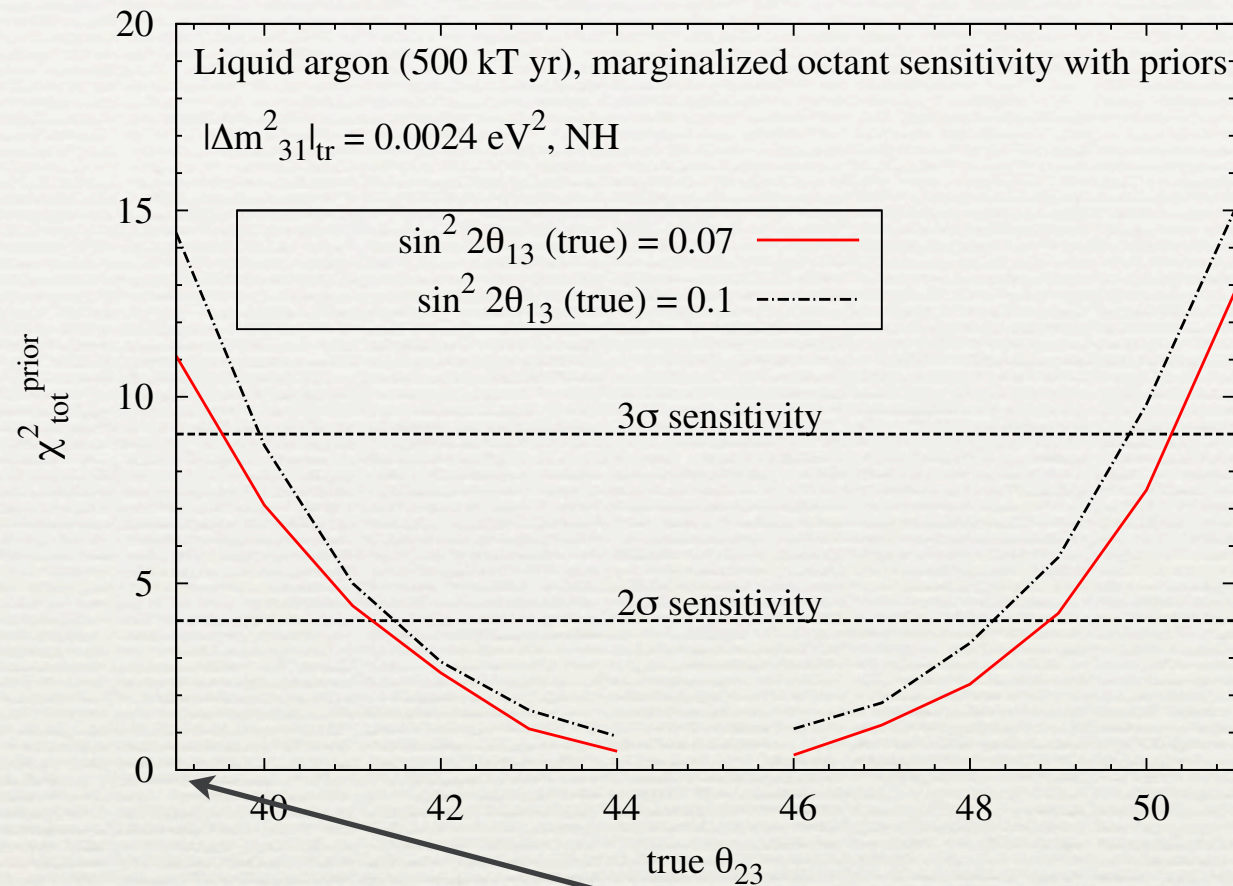


If θ_{13} in 1 sigma range of RENO/DB then 250 kT-yr exposure gives ~ 5 sigma sensitivity for NH & IH

If present error in θ_{13} is reduced by half, a 100 kT-yr exposure is capable of delivering a 4 sigma determination of the hierarchy.



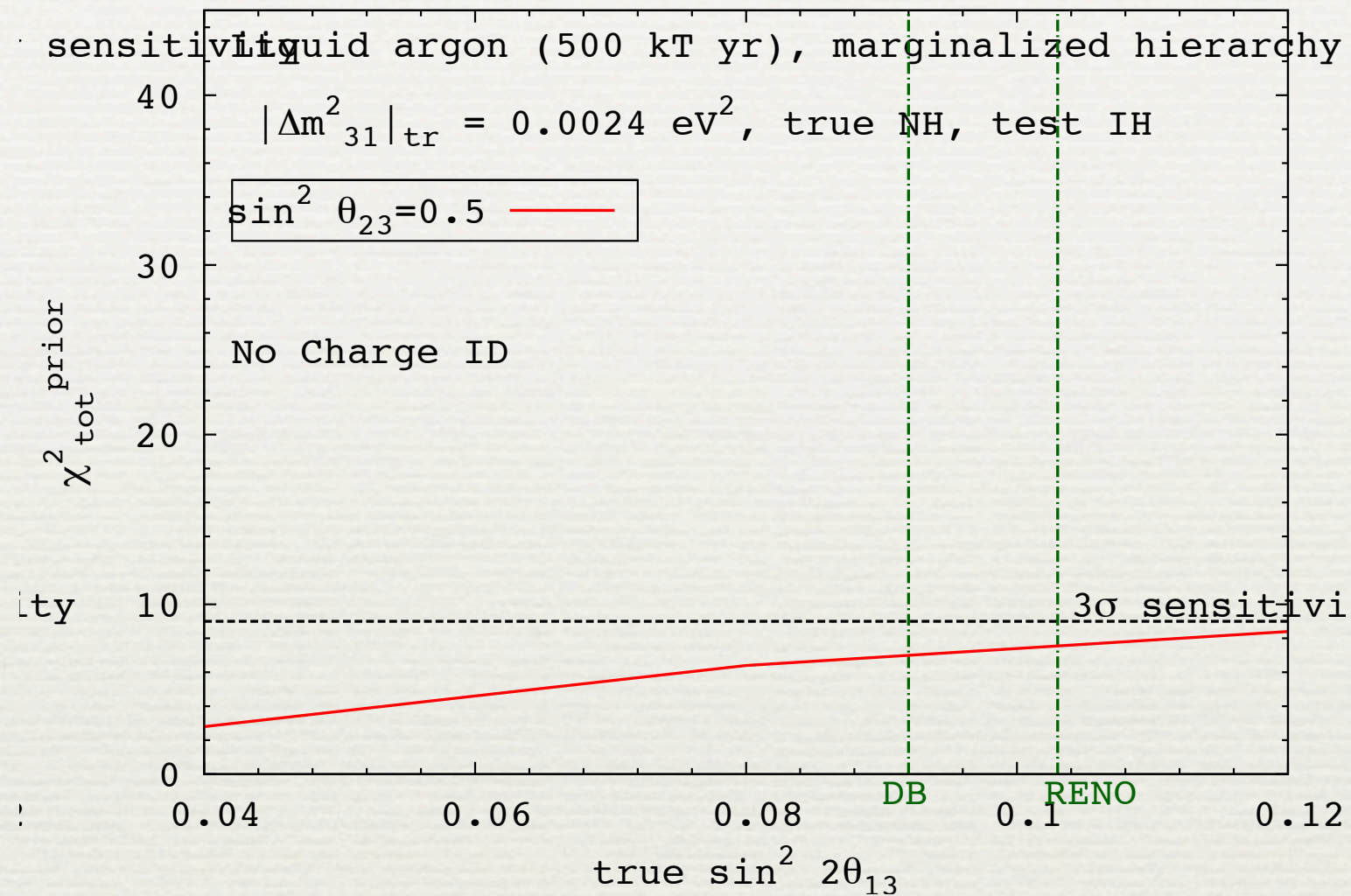
Results.....Octant



Present global best fit
 Fogli et al arXiv 1205.5254

Octant sensitivity due to $\sin^4 \theta_{23}$ term in survival probability

>2 sigma determination for 500 kT-yr if $|\theta_{23} - \pi/4| > 3.5^\circ$



3 sigma hierarchy determination with about 600 kT-yr exposure if no charge id.

Conclusions.....

Of three important goals in neutrino physics, two (hierarchy and octant) may be reachable significantly sooner than the third (CP).

A magnetized LAr-TpC, given its superior spatial, charge-id and calorimetric properties, can determine the hierarchy to 5 sigma with 250-kt-yr at current errors of θ_{13}

If this error is halved by future measurements, it can deliver a 4 sigma determination with 100 kT-yr exposure.

It can determine the octant to ~3 sigma with 500 kT-yr if θ_{23} at current global best fit value.

Conclusions.....

A 10-20 kT LAr-TPC with or without a magnetic field, sited underground to detect beam as well as atmospheric neutrinos would deliver even superior physics results. (work in progress.....)

Thank you for your attention!